

# Did DESI DR2 truly reveal dynamical dark energy?

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A fundamental question in cosmology is whether dark energy evolves over time, a topic that has gained prominence since the discovery of cosmic acceleration. Recently, the DESI collaboration has reported increasing evidence for evolving dark energy using combinations of cosmic microwave background (CMB), type Ia supernova (SN), and their new measurements of baryon acoustic oscillations (BAO). However, our analysis reveals that these combinations are problematic due to clear tensions among the CMB, BAO and SN datasets. Consequently, DESI’s claim of dynamical dark energy (DDE) is not robust. A more reliable approach involves constraining the evolution of dark energy using each dataset independently. Through a statistical comparison for each dataset, on average, we find that DDE is strongly preferred over the  $\Lambda$ CDM model. This suggests that DDE likely exists, although its real parameter space remains elusive due to weak constraints on the dark energy equation of state and inconsistencies among the datasets. Interestingly, when considering DDE, none of the individual datasets—including CMB, DESI DR2, Pantheon+, Union3, and DESY5—can independently detect cosmic acceleration at a significant level. Our findings not only clarify the current understanding of the nature of dark energy but also challenge the established discovery of cosmic acceleration and the long-held notion that dark energy exerts negative pressure. Both individual and combined datasets suggest that the ultimate fate of the universe is likely to be dominated by matter rather than dark energy.

*Introduction.* The standard cosmological model,  $\Lambda$  cold dark matter ( $\Lambda$ CDM), which has been confirmed by the CMB [1–5], BAO [6–12], and SN observations [13, 14], can successfully characterize various kinds of physical phenomena such as the cosmic acceleration on large scales and the clustering of matter on small scales [15]. However, it confronts two intractable challenges, namely the cosmological constant conundrum [16, 17] and the coincidence problem, while suffering from the emergent cosmic tensions such as the so-called Hubble constant ( $H_0$ ) tension and the matter fluctuation amplitude ( $S_8$ ) discrepancy [18–22]. It is very logically reasonable to query the validity of  $\Lambda$ CDM in verifying the fundamental theory and depicting the background dynamics and structure formation of the universe. So far, to solve these discrepancies, there are a great deal of alternative scenarios proposed by different authors (see [21, 22] for reviews). It is worth noting that, besides the theoretical developments, more importantly, we require new independent and powerful probes with higher precision to give definite answers on some core puzzles. To achieve this goal, a very promising probe is BAO.

BAO are regular and periodic matter density fluctuations of the universe [15], which originate from sound waves induced by hot baryon-photon plasma before the epoch of recombination. The characteristic scale of BAO, approximately 150 Mpc, which is the maximum distance that the acoustic waves could travel in the primordial plasma before the plasma cooled to the point where it became neutral atoms at the recombination epoch, serves as a standard ruler in cosmology. Many BAO experiments such as 2dF [7], 6dF [8], SDSS [6] and eBOSS [11, 12]

map the late-time expansion history of the universe by measuring the apparent size of this ruler at different redshifts. BAO are very clean probes to explore the evolution of the universe over time, which is unaffected by the nonlinear physics on small scales and relatively robust to systematic uncertainties compared to other cosmological probes.

Recently, the DESI collaboration give the substantial evidence of DDE [23], based on their measurements of BAO in galaxy, quasar and Lyman- $\alpha$  forest tracers from the first data release (DR1) of the Dark Energy Spectroscopic Instrument (DESI) [24, 25]. Interestingly, this DDE evidence is enhanced [26, 27] by the DESI’s second data release (DR2) including more than 14 million galaxies and quasars, based on three years of operation [28].

In theory, DDE predicts: (i) the equation of state (EoS) and energy densities of dark energy (DE) evolves over time; (ii) different expansion history of the universe from  $\Lambda$ CDM; (iii) the rate at which cosmic structures like galaxies and clusters form can be changed; (iv) the fate of the universe can be significantly affected. If DDE is finally demonstrated to be true, it will indicate that the vacuum is not empty and it does have matter. Therefore, the DESI’s finding of DDE evidence is crucial for theory. Up to date, the addition of CMB and SN data to DESI DR2 leads to  $2.8 - 4.2\sigma$  evidence of DDE [26], depending on which SN sample is used. Although the unprecedented precision and number of data points lead to the evidence of DDE, we should be very cautious about these results. The key reason is that the DDE evidence is derived by the data combination of CMB, DESI DR2

and SN, not from each probe independently. If each of CMB, DESI DR2 and SN gives consistent constraints on the cosmological parameters including the matter fraction and DE EoS, one can safely claim that DDE do exist in the late-time universe. Based on this concern, one has to seriously question: whether does DESI truly see DDE? Our results show that it is too early to claim the existence of DDE using the combinations of CMB, DESI DR2 and SN data, but independent datasets still give strong statistical preferences of DDE over  $\Lambda$ CDM. Future high precision observations can help clarify the status of DE.

*Basics.* In the theory of general relativity [29], considering a homogeneous and isotropic universe, the Friedmann equations read as  $H^2 = (8\pi G\rho)/3$  and  $\ddot{a}/a = -4\pi G(\rho + 3p)/3$ , where  $a$  is the scale factor,  $H$  is the cosmic expansion rate and  $\rho$  and  $p$  are the mean energy density and pressure of different species including radiation, baryons, dark matter and DE. Combining two Friedmann equations, one can express the dimensionless Hubble parameter  $E(z) \equiv H(z)/H_0$  for a flat Chevallier-Polarski-Linder (CPL) universe [30, 31] as

$$E(a) = \left[ \Omega_m a^{-3} + (1 - \Omega_m) a^{-3(1+\omega_0+\omega_a)} e^{3\omega_a(a-1)} \right]^{\frac{1}{2}}, \quad (1)$$

where  $\Omega_m$  is the matter fraction. It reduces to  $\Lambda$ CDM when  $\omega_0 = -1$  and  $\omega_a = 0$ .

*Data and methodology.* We use the Planck 2018 high- $\ell$  `plik` temperature (TT) likelihood at multipoles  $30 \leq \ell \leq 2508$ , polarization (EE) and their cross-correlation (TE) data at  $30 \leq \ell \leq 1996$ , and the low- $\ell$  TT `Commander` and `SimA11` EE likelihoods at  $2 \leq \ell \leq 29$  [32]. We adopt conservatively the Planck lensing likelihood [33] from `SMICA` maps at  $8 \leq \ell \leq 400$ . We use the most recent 13 BAO measurements from DESI DR2 including the BGS, LRG1, LRG2, LRG3+ELG1, ELG2, QSO and Ly $\alpha$  samples at effective redshifts  $z_{\text{eff}} = 0.295, 0.51, 0.706, 0.934, 1.321, 1.484$  and  $2.33$ , respectively [26–28]. To completely explore the DE EoS at late times, we adopt three well-calibrated SN compilations: (i) Pantheon+ consisting of 1701 data points from 18 different surveys in  $z \in [0.00122, 2.26137]$  [34]; (ii) Union3 with 22 spline-interpolated data points derived by 2087 SN from 24 different surveys in  $z \in [0.05, 2.26]$  [35]; (iii) DESY5 including 1735 effective data points in  $z \in [0.025, 1.130]$  [36].

To calculate the background dynamics of the universe and theoretical power spectra, we use the Boltzmann solver `CAMB` [37]. To implement the Bayesian analysis, we employ the Monte Carlo Markov Chain (MCMC) method to infer the posterior distributions of model parameters using the publicly available package `Cobaya` [38]. We assess the convergence of MCMC chains via the Gelman-Rubin criterion  $R-1 \lesssim 0.01$  [39] and analyze them using `Getdist` [40].

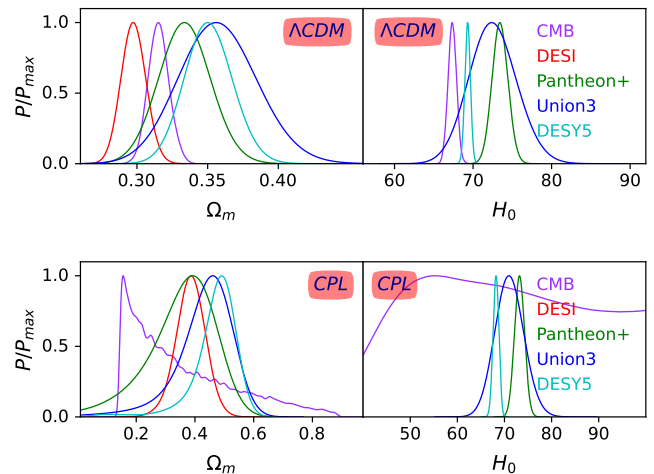


FIG. 1: One-dimensional posterior distributions of the parameters  $\Omega_m$  and  $H_0$  from CMB, DESI DR2 and SN observations in the  $\Lambda$ CDM (upper) and CPL (lower) models, respectively.

We use the following uniform priors for free parameters: the baryon fraction  $\Omega_b h^2 \in [0.005, 0.1]$ , cold dark matter fraction  $\Omega_c h^2 \in [0.001, 0.99]$ , acoustic angular scale at the recombination epoch  $100\theta_{MC} \in [0.5, 10]$ , scalar spectral index  $n_s \in [0.8, 1.2]$ , amplitude of the primordial scalar power spectrum  $\ln(10^{10} A_s) \in [2, 4]$ , optical depth  $\tau \in [0.01, 0.8]$ , present-day DE EoS  $\omega_0 \in [-15, 20]$  and the amplitude of DE evolution  $\omega_a \in [-30, 10]$ . To produce a matter-dominated era at high redshifts, we impose the condition  $\omega_0 + \omega_a < 0$  in the Bayesian analysis. The reason why we take such wide priors for the parameter pair  $(\omega_0, \omega_a)$  is that a large enough parameter space can completely present the constraining power of the DESI DR2 BAO measurements [41].

*$\Omega_m$  and  $H_0$  tensions from independent probes.* The DESI collaboration have noticed that there is a  $\Omega_m$  tension among CMB, DESI DR2 and three SN samples under  $\Lambda$ CDM [26]. SN datasets clearly prefer a larger matter fraction than CMB and DESI DR2. However, the accompanying  $H_0$  tension is not reported. In Fig.1, we find that Pantheon+, Union3 and DESY5 exhibit  $H_0$  tensions with CMB [1] at  $5.31\sigma$ ,  $1.69\sigma$  and  $3.03\sigma$  levels, respectively. Although three SN samples are well calibrated, they suffer from internal inconsistencies in their preferred  $H_0$  values. For instance, Pantheon+ is in a  $3.86\sigma$  tension with DESY5. The main goal of DESI is exploring the nature of DE including its possible dynamics. Based on the new motivation that DDE could help resolve the  $\Omega_m$  tension in  $\Lambda$ CDM, we implement constraints on the CPL DDE and find that CMB prefers a smaller  $\Omega_m$ , while DESI DR2 and SN prefer larger matter fractions. Interestingly,  $\Omega_m$  tensions between DESI DR2 and CMB, Pantheon+, Union3, and DESY5 are well alleviated from

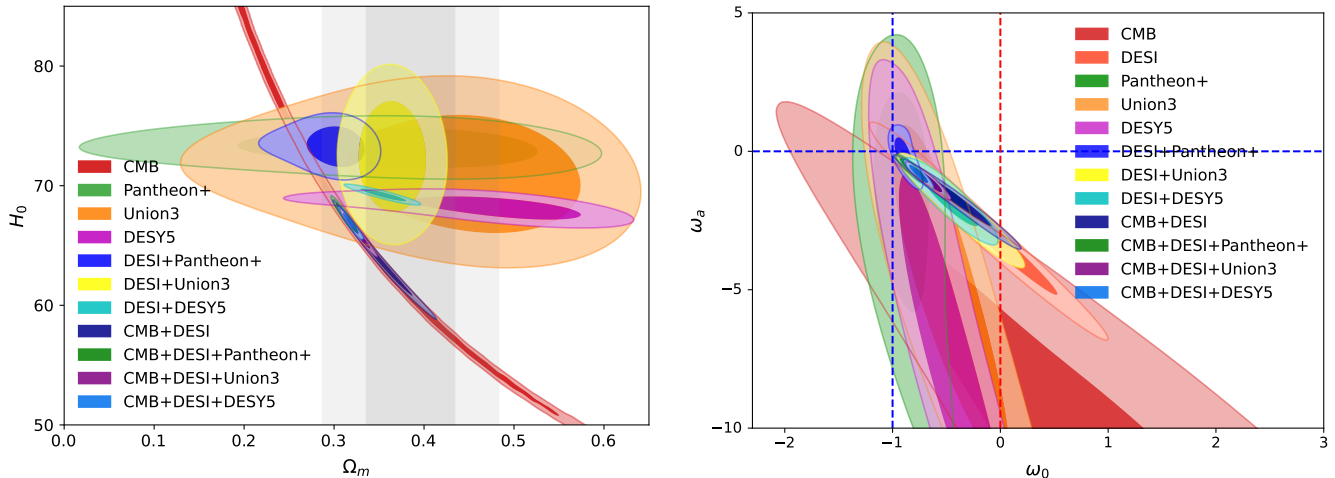


FIG. 2: Two-dimensional posterior distributions of the parameter pairs  $(\Omega_m, H_0)$  and  $(\omega_0, \omega_a)$  from different datasets in the CPL model. The vertical shaded grey regions are the constrained  $1\sigma$  and  $2\sigma$   $\Omega_m$  values from DESI DR2. The cross point of blue dashed lines represents  $\Lambda$ CDM. The red dashed line denotes  $\omega_0 = 0$ .

$1.57\sigma$ ,  $1.83\sigma$ ,  $2.10\sigma$  and  $2.81\sigma$  to  $0.14\sigma$ ,  $0.21\sigma$ ,  $0.49\sigma$  and  $0.94\sigma$ , respectively. It seems that DDE solves the  $\Omega_m$  discrepancies, however, the price is DESI DR2 and SN give larger  $\Omega_m$  (e.g.  $\sim 0.5$ ) in CPL than in  $\Lambda$ CDM, while the  $H_0$  tension between Pantheon+ and DESY5 increases from  $3.86\sigma$  to  $4.29\sigma$  (see Tab.I). Therefore, the well-known fact that the late universe is dominated by DE is largely challenged in CPL. Note that CMB cannot constrain  $H_0$  in CPL, due to the weak constraint on the DE EoS.

*Inconsistencies from data combinations.* The strong evidence for DDE primarily comes from the combination of CMB, BAO, and SN data. However, this result accommodates the inconsistencies among them. Hence, it is reasonable to question the validity of the combination. In Fig.2, the insensitivity of CMB to  $(\omega_0, \omega_a)$  leads to a much weaker constraint on  $H_0$  and  $\Omega_m$  in CPL than that in  $\Lambda$ CDM. Fortunately, the strong anti-correlation between  $H_0$  and  $\Omega_m$  remains. Combining CMB with BAO and SN, this anti-correlation will largely help compress the parameter space and consequently tighten constraints on  $(\omega_0, \omega_a)$ . Unfortunately, there are beyond  $1\sigma$  tensions between CMB and SN. Especially, DESY5 gives a  $\sim 2\sigma$  tension with CMB. Even worse, the addition of DESI DR2 to SN leads to beyond  $2\sigma$  tensions with CMB, particularly, DESY5 plus DESI DR2 gives a beyond  $3\sigma$  tension with CMB. These tensions can also be captured in  $r_d$  values. Adding DESI DR2 to Pantheon+, Union3 and DESY5 leads to a  $5.30\sigma$ ,  $2.86\sigma$  and  $4.26\sigma$  tension with CMB when estimating  $r_d$ , respectively. Notice that there is a beyond  $2\sigma$  tension between Pantheon+ plus DESI DR2 and DESY5 plus DESI DR2. This suggests a clear discrepancy among SN samples. Interestingly, CMB plus DESI DR2 is in beyond  $3\sigma$ , beyond  $2\sigma$  and  $\sim 1.5\sigma$  ten-

sions with Pantheon+, DESY5 and Union3, respectively. All these discrepancies demonstrate that the DDE evidences from combined constraints are problematic.

Furthermore, we give more complete constraints on  $(\omega_0, \omega_a)$  using a wider prior than that the Planck and DESI collaborations used [1, 26, 27]. We find that CMB is in a  $\sim 2\sigma$  tension with DESI DR2 and three SN samples are basically inconsistent with CMB at  $\sim 1\sigma$  level. Since DESI DR2 with a clear different degeneracy direction of  $(\omega_0, \omega_a)$  from SN has a  $\sim 1\sigma$  tension with SN, the DDE evidences from DESI DR2 plus CMB or SN are problematic. Additionally, even for the SN-only case, there is also an inconsistency among SN samples, i.e., Pantheon+ is consistent with  $\Lambda$ CDM within  $1\sigma$  level, while Union3 and DESY5 exhibit beyond  $1\sigma$  tensions with  $\Lambda$ CDM. Therefore, we should not place too much trust in the constraints derived from any pairwise combination of the datasets. Starting from this viewpoint, the reason why  $\Omega_m$  values from the combination of CMB, DESI DR2 and SN are so close to  $\Omega_m = 0.3153 \pm 0.0073$  [1] from the CMB-only constraint on  $\Lambda$ CDM should be a coincidence.

*Solution.* Due to the unprecedented precision and increasing number of data points, we should be cautious when studying possible new physics using the data combinations, which may bias the results away from the truth. Although there are discrepancies among CMB, BAO and SN data, an undeniable fact is that they all prefer the region of  $\omega_0 > -1$  and  $\omega_a < 0$ . Especially, DESI DR2, giving a strong enough anti-correlation of  $(\omega_0, \omega_a)$ , shows a  $\sim 2\sigma$  preference with a relatively high precision. Concerning this, we implement a statistical comparison between CPL and  $\Lambda$ CDM. For CMB, we find that the Bayesian factor  $\ln B_{ij} = 6.34 > 5$  [45], indicating a

TABLE I: Mean values and  $1\sigma$  (68%) uncertainties of free parameters from different datasets in the CPL model. We quote the  $2\sigma$  (95%) upper limit of  $\omega_0$  in the CMB-only case. The symbols “★” and “◆” denote unconstrained parameters by data and parameters with poor constraints, respectively. Here  $h \equiv H_0/100$  in units of  $\text{km s}^{-1} \text{Mpc}^{-1}$ .

Parameter	$\omega_0$	$\omega_a$	$\Omega_m$	$H_0$	$r_d$	$hr_d$
CMB	$2.4 \pm 1.9$	$< -3.62$	$0.344^{+0.054}_{-0.200}$	◆	$147.32 \pm 0.27$	$103^{+20}_{-30}$
DESI	$-0.17 \pm 0.44$	$-2.8 \pm 1.6$	$0.385 \pm 0.049$	★	★	$91.5^{+4.4}_{-4.9}$
Pantheon+	$-0.89 \pm 0.17$	$-2.1^{+3.2}_{-1.8}$	$0.360^{+0.130}_{-0.086}$	$73.2 \pm 1.0$	★	★
Union3	$-0.45^{+0.28}_{-0.40}$	$-5.4^{+4.7}_{-3.1}$	$0.437^{+0.100}_{-0.066}$	$71.1 \pm 3.0$	★	★
DESY5	$-0.35^{+0.30}_{-0.41}$	$-9.0^{+5.4}_{-4.5}$	$0.471^{+0.075}_{-0.043}$	$68.20 \pm 0.60$	★	★
DESI+Pantheon+	$-0.885 \pm 0.061$	$-0.19 \pm 0.46$	$0.299^{+0.025}_{-0.016}$	$73.3 \pm 1.0$	$136.1 \pm 2.1$	$99.72 \pm 0.93$
DESI+Union3	$-0.37 \pm 0.23$	$-2.07 \pm 0.82$	$0.365 \pm 0.024$	$72.4 \pm 3.0$	$129.3 \pm 6.3$	$93.5 \pm 2.6$
DESI+DESY5	$-0.47 \pm 0.17$	$-1.75 \pm 0.61$	$0.354 \pm 0.016$	$69.24 \pm 0.34$	$136.6 \pm 2.5$	$94.6 \pm 2.1$
CMB+Pantheon+	$-0.868 \pm 0.092$	$-0.56^{+0.48}_{-0.42}$	$0.311^{+0.011}_{-0.013}$	$67.8 \pm 1.2$	$147.13 \pm 0.25$	$99.7 \pm 1.8$
CMB+Union3	$-0.63 \pm 0.14$	$-1.39^{+0.68}_{-0.60}$	$0.319^{+0.012}_{-0.014}$	$66.9 \pm 1.3$	$147.17 \pm 0.25$	$98.5 \pm 2.0$
CMB+DESY5	$-0.73 \pm 0.10$	$-1.04 \pm 0.51$	$0.3151^{+0.0094}_{-0.0110}$	$67.3^{+1.10}_{-0.95}$	$147.15 \pm 0.25$	$99.1^{+1.6}_{-1.4}$
CMB+DESI	$-0.41^{+0.21}_{-0.25}$	$-1.74^{+0.75}_{-0.59}$	$0.353^{+0.022}_{-0.025}$	$63.7 \pm 2.0$	$147.16 \pm 0.23$	$93.7 \pm 3.1$
CMB+DESI+Pantheon+	$-0.843 \pm 0.054$	$-0.58^{+0.23}_{-0.19}$	$0.3108 \pm 0.0058$	$67.62 \pm 0.60$	$147.28 \pm 0.22$	$99.59 \pm 0.89$
CMB+DESI+Union3	$-0.673 \pm 0.087$	$-1.04^{+0.30}_{-0.27}$	$0.3268 \pm 0.0086$	$66.02 \pm 0.84$	$147.22 \pm 0.21$	$97.2 \pm 1.2$
CMB+DESI+DESY5	$-0.757 \pm 0.057$	$-0.82 \pm 0.22$	$0.3184 \pm 0.0057$	$66.86 \pm 0.57$	$147.24 \pm 0.22$	$98.44 \pm 0.84$

strong evidence of CPL from CMB. For DESI DR2, Pantheon+, Union3 and DESY5, we obtain  $\Delta\text{BIC} = 2.25, 25.63, 3.92$  and  $11.49$  [46], respectively, indicating positive evidences from DESI DR2 and Union3 and strong evidences from Pantheon+ and DESY5 for DDE. Hence, DDE likely exists, however, we cannot determine the real parameter space of DDE according to such independent probes, due to their weak constraints on  $(\omega_0, \omega_a)$  and inconsistencies. If future observations (CMB, BAO or SN) with increasing precision can provide inconsistent constraints on DDE, we can safely claim its existence.

*The fate of the universe.* The existence of DDE will affect the composition, fate, expansion history and structure formation of the universe. The fact that three SN datasets allow  $\Omega_m > 0.5$  in CPL means that the late universe could be matter dominated not DE dominated. At least, CMB, DESI DR2 and SN all independently allow a large  $\Omega_m \sim 0.5$ . Interestingly, CMB, DESI DR2, Union3 and DESY5 do not rule out  $\omega_0 > 0$  and  $\omega_0 > 1/[3(\Omega_m - 1)]$  [47]. Only Pantheon+ finds a beyond  $1\sigma$  hint of late-time cosmic acceleration. One has to question whether the universe is accelerating now. CMB, DESI DR2, Union3 and DESY5 give, respectively,  $\omega_0 = 2.4 \pm 1.9, -0.17 \pm 0.44, -0.45^{+0.28}_{-0.40}$  and  $-0.35^{+0.30}_{-0.41}$ , allowing that current universe could be slowing down or moving at a constant speed. Particularly, CMB gives a  $1.3\sigma$  hint of  $\omega_0 > 0$ , allowing that DE has a positive pres-

sure. Our findings not only profoundly challenge the understanding of cosmic acceleration, but also challenge the long-held notion that the pressure of DE is negative. Interestingly, we find that both independent and combined datasets prefer that the ultimate destiny of the universe is completely filled with matter not DE, i.e.,  $\Omega_m = 1$ . At some point in the distant future, matter will dominate the evolution of the universe. Finally, the universe will miraculously stop moving [48]. More details about the fate of the universe will be shown in [49].

*In summary,* the DESI DR2 data, when combined with CMB and SN data, appears to provide increasing evidence for Dynamical Dark Energy (DDE). However, our analysis demonstrates that this conclusion is not robust. By constraining the CPL model with independent observations and various data combinations, we identify significant tensions between different datasets. These tensions lead to problematic constraints when incompatible datasets are combined.

Although the combined constraints yield  $\Omega_m$  and  $H_0$  values close to those from the CMB-only constraint on  $\Lambda\text{CDM}$ , we believe this to be coincidental. In the CPL model, the effect of the cosmological constant in  $\Lambda\text{CDM}$  is partially replaced by a larger matter fraction and DDE.

The combination of CMB, DESI DR2, and SN data results in a lower  $H_0$  value because DESI DR2 indicates a high  $\Omega_m = 0.385 \pm 0.049$ . This high matter density

directly leads to a low  $H_0 = 63.7 \pm 2.0 \text{ km s}^{-1} \text{ Mpc}^{-1}$  when combined with CMB data, due to the strong anti-correlation between  $H_0$  and  $\Omega_m$  in CMB observations. Given the crucial role of  $H_0$  in the universe's background evolution, we explore the effects of  $H_0$  and  $H_0 r_d$  on the parameters  $\omega_0$  and  $\omega_a$  for DESI DR2 in the supplementary material (SM). Additionally, we find that allowing for a free lensing amplitude  $A_L$  can reduce the significance of DDE by approximately  $1\sigma$ , as  $A_L$  and DDE are degenerate in the CMB lensing potential (see SM).

It is important to note that as the precision and volume of cosmological data increase, so do the tensions between different probes. This growing statistical complexity challenges our understanding of the data, underlying physics, and potential systematics. We are currently at a critical juncture in addressing these challenges.

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- [45] For CMB, choosing  $\Lambda$ CDM as the reference model, we use the `MCEvidence` package [44] to calculate the Bayesian evidences of CPL and  $\Lambda$ CDM,  $\varepsilon_i$ , and the Bayes factor,  $B_{ij} = \varepsilon_i/\varepsilon_j$ , where  $\varepsilon_j$  is the statistical evidence of the reference model. We use a revised and more conservative version of the so-called Jeffreys’ scale [42], i.e.,  $\ln B_{ij} = 0-1$ ,  $\ln B_{ij} = 1-2.5$ ,  $\ln B_{ij} = 2.5-5$  and  $\ln B_{ij} > 5$  implying an *inconclusive*, *weak*, *moderate* and *strong* preference of the model  $i$  over the reference model  $j$ , respectively. Note that data prefers the reference model if  $\ln B_{ij} < 0$  for an experiment. Here the reference model is  $\Lambda$ CDM.
- [46] For DESI DR2 and SN, we implement the statistical comparison using the Bayesian information criterion (BIC) [43], which is defined as  $\text{BIC} = \chi_{\min}^2 + n \ln N$ , where  $n$  and  $N$  denotes the numbers of model parameters and data points. It is noteworthy that the relative values of this criterion between different models are more useful than the absolute values to distinguish which model is preferred by observations. Hence, we define the relative variation  $\Delta\text{BIC}$  between models as follows  $\Delta\text{BIC} = \text{BIC}^{\text{model}} - \text{BIC}^{\Lambda\text{CDM}}$ , where we have chosen the  $\Lambda$ CDM model as the reference model. Generally, the differences  $\Delta\text{BIC} \sim 2$  and  $\Delta\text{BIC} \sim 6$  indicate a positive and strong evidence against the reference model, respectively.
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## Supplementary Material for “Did DESI DR2 truly reveal dynamical dark energy?”

In this supplementary material, first of all, we study the effect of CMB lensing amplitude  $A_L$  on the DDE constraints using the combinations of CMB, DESI DR2 and SN observations. Since the DE EoS is closely related to the present-day cosmic expansion rate, we then investigate the effect of  $H_0$  and  $H_0 r_d$  on the DDE constraints using only DESI DR2 BAO measurements. Finally, we compare the CMB-only constraints on the CPL DDE and  $\Lambda$ CDM.

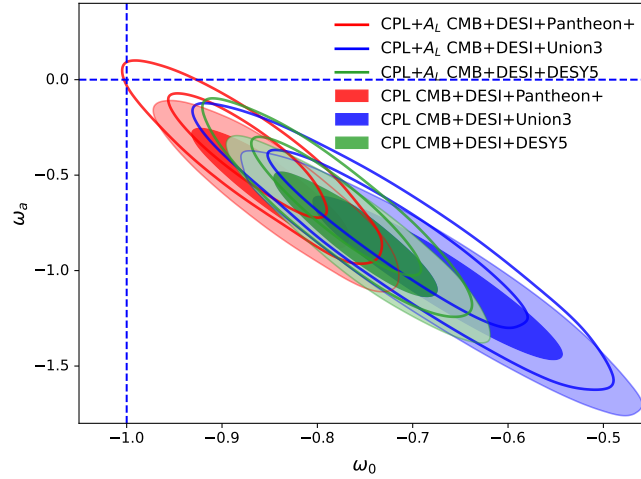


FIG. 1: Two-dimensional posterior distributions of the parameter pairs  $(\omega_0, \omega_a)$  from CMB, DESI DR2 and SN observations in the CPL model with and without the lensing amplitude  $A_L$ . The cross point of blue dashed lines corresponds to  $\Lambda$ CDM.

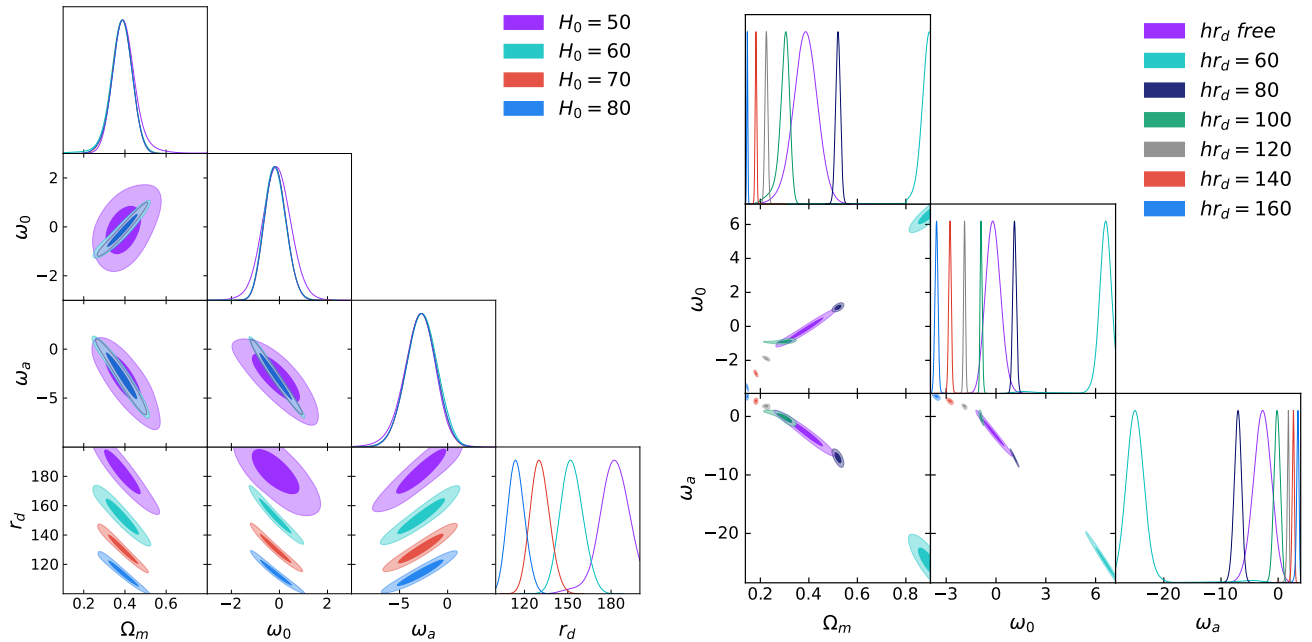


FIG. 2: One-dimensional and two-dimensional posterior distributions of model parameters in the CPL model from the DESI DR2 data when considering different values of  $H_0$  and  $hr_d$ , respectively. For the case of  $hr_d$ , we use a free  $hr_d$  as a comparison.



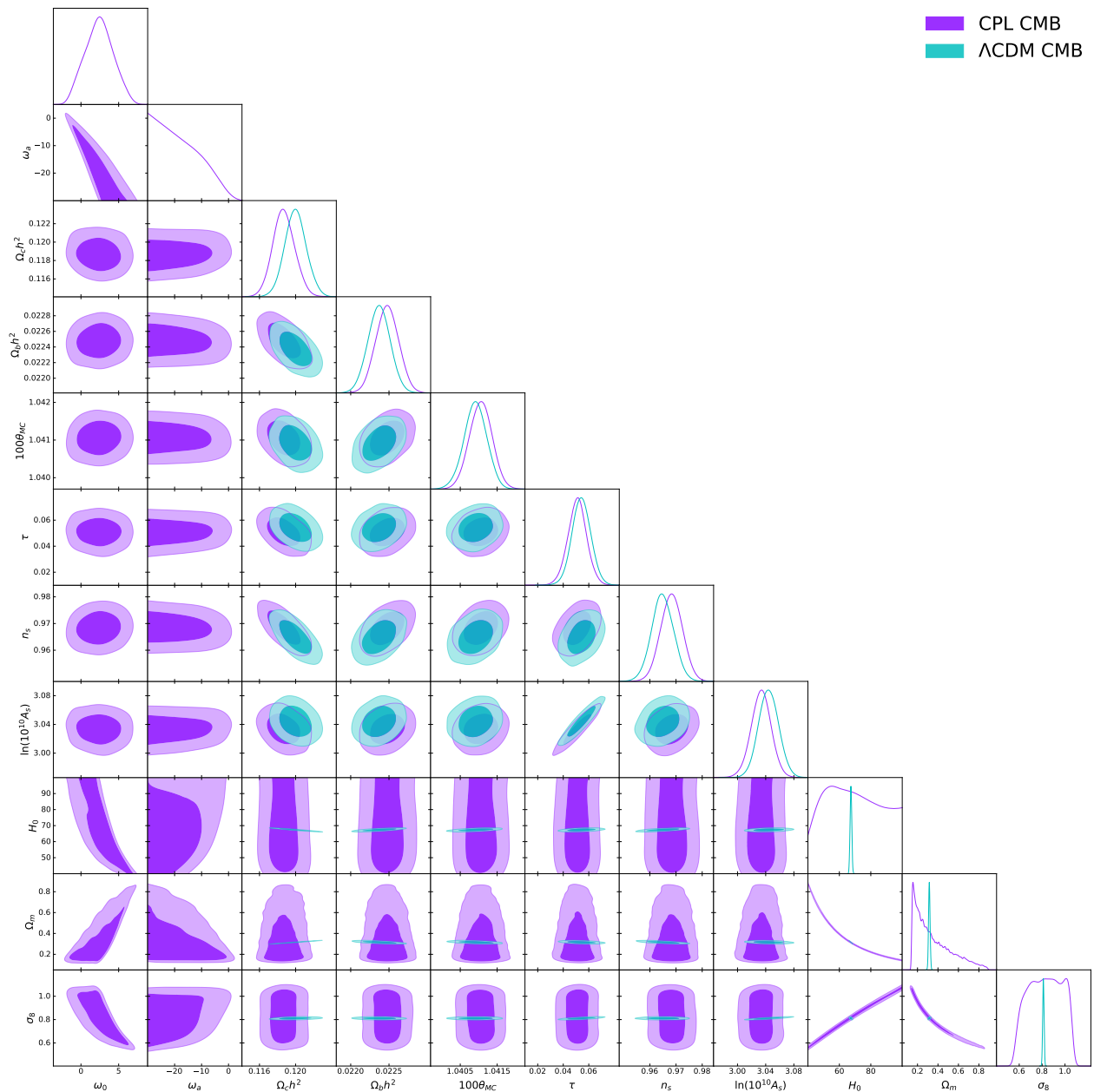


FIG. 3: One-dimensional and two-dimensional posterior distributions of model parameters in the CPL and  $\Lambda$ CDM models from the CMB observations.

### A. The effect of lensing amplitude on DDE

As we know, the CMB lensing anomaly is strongly degenerated with many new physics on cosmic scales, e.g., modified gravity [1, 2], cosmic curvature [3] and effective field theory of DE [1]. Since both DDE and  $A_L$  have impacts on the gravitational potential, considering  $A_L$  could change the constraints on  $(\omega_0, \omega_a)$ . Employing the data combination of CMB, DESI DR2 and SN, we find that a free  $A_L$  can reduce a  $\sim 1\sigma$  significance of DDE (see Fig.1). This result holds for all three SN samples considered here.



### B. The effects of $H_0$ and $hr_d$ on DDE

In theory, BAO require the information of the comoving sound horizon  $r_d$  at the drag epoch from CMB [ ] to help determine  $H_0$ , since BAO cannot independently constrain  $H_0$ . However, it is interesting to see how the correlations between  $r_d$  and the DE EoS (or matter density ratio) vary over different  $H_0$  values. In Fig.2, as expected, varying  $H_0$  hardly changes the constraints on the matter fraction. Larger  $H_0$  leads to smaller  $r_d$ .  $r_d$  inherits the anti-correlations between  $H_0$  and  $\Omega_m$  (or  $\omega_0$ ) and the positive correlation between  $H_0$  and  $\omega_a$ . It is worth noting that the whole parameter space will be enlarged when taking small  $H_0$  values. Furthermore, we study how the compound parameter  $H_0r_d$  impacts the DDE constraints. Similar to the case of  $H_0$ , larger  $H_0r_d$  gives smaller  $\Omega_m$  and  $\omega_0$  as well as larger  $\omega_a$ . The constraints from the case of free  $H_0r_d$  fall well in between those from  $H_0r_d = 80$  and  $100$ , because it gives  $H_0r_d = 91.5^{+4.4}_{-4.9}$  (see the table in the main text).

### C. CPL vs $\Lambda$ CDM in light of CMB

In Fig.3, we make a comparison between the CPL DDE and  $\Lambda$ CDM models. Overall, DDE gives similar constraints on six basic parameters, although there are small shifts in some parameters. It is noteworthy that CMB is phenomenologically insensitive to the late-time DE EoS. This insensitivity leads to a poor constraint on  $(\omega_0, \omega_a)$ . Therefore, one cannot well constrain the background quantities such as  $H_0$  and  $\Omega_m$  in the CPL DDE model.

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